

Magnetic Domain Patterns Depending on the Sweeping Rate of Magnetic Fields

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The domain patterns in a thin ferromagnetic film are investigated in both experiments and numerical simulations. Magnetic domain patterns under a zero field are usually observed after an external magnetic field is removed. It is demonstrated that the characteristics of the domain patterns depend on the decreasing rate of the external field, although it can also depend on other factors. Our numerical simulations and experiments show the following properties of domain patterns: a sea-island structure appears when the field decreases rapidly from the saturating field to the zero field, while a labyrinth structure is observed for a slowly decreasing field. The mechanism of the dependence on the field sweeping rate is discussed in terms of the concepts of crystallization.

KEYWORDS: domain pattern, pattern formation, garnet thin film, domain dynamics, time-dependent field

There are a large number of physical and chemical systems that display domain patterns.¹ Ferromagnetic garnet thin films also show domain patterns, which can have various kinds of structures. For example, when an external field is applied, a hexagonal lattice, bubbles, or stripes are observed. Under a zero field, usually a labyrinth structure is observed, although a parallel-stripe structure is more stable than a labyrinth structure. In fact, it is observed in experiments that a labyrinth structure changes to a parallel-stripe structure under field cycles.^{2,3} However, the domain patterns observed under a zero field do not always have a labyrinth or parallel-stripe structure. In some samples, under a certain condition, a sea-island structure is observed, which consists of many small up-spin (down-spin) domains surrounded by a “sea” of down (up) spins.

In order to observe the magnetic domain patterns under a zero field in experiments, usually an external magnetic field is applied once up to the saturation field, and then it is removed. Here, we will focus on the domain patterns under a zero field and their dependence on the decreasing rate of the external field, although it is natural that the characteristics of the observed patterns can depend on the properties of the samples, *e.g.*, disorder, anisotropy, interactions among spins, and so on. In fact, there have been several simulations that show how domain patterns change depending on such properties under a very slowly changing field.^{4–6} There also exist some works on the dependence of hysteresis on the frequency and amplitude of an oscillating field.^{7–9} However, for the characteristics of domain patterns under a zero field, their dependence on the changing rate of a time-dependent field has not been investigated so far.

In this letter, we report both the experimental observations and the numerical simulations of domain patterns. In particular, we show how the characteristics of domain patterns depend on the sweeping rate of the magnetic field. First, the experimental procedure and results are explained. Next, the model of the system and numeri-

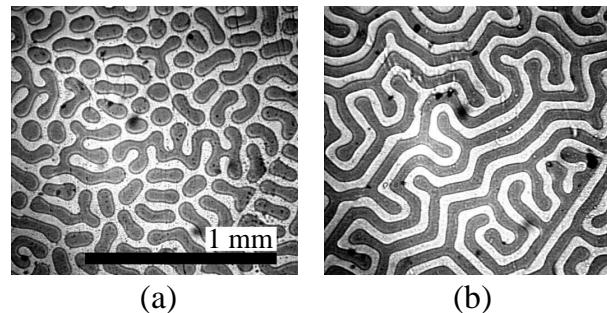


Fig. 1. Experimentally observed domain patterns under a zero field. The external field was decreased at a (a) fast rate and (b) slow rate. (See the text.)

cal procedure are described. After displaying numerically simulated domain patterns, we study the field dependence of the characteristics of domain patterns, *i.e.*, the number of domains and the domain area. The mechanism of the dependence will also be discussed.

Experiments were performed on ferrimagnetic garnet films of composition $(\text{BiGdY})_3(\text{FeGa})_5\text{O}_{12}$ grown epitaxially on the (111) face of gadolinium gallium garnet substrates. This film has strong uniaxial magnetic anisotropy with the easy axis perpendicular to the film. The perpendicular saturating field is about 120 Oe. Magnetic domain structures are easily observed by using a polarizing microscope with the help of the magneto-optical Faraday effect. Figure 1(a) shows the domain pattern under a zero field when the magnetic field has been decreased from 220 to 0 Oe at a fast sweeping rate of 2×10^5 Oe/s. This picture shows a sea-island structure. The observed structure at a sweeping rate of 10 Oe/s is shown in Fig. 1(b). At this slow rate, the domain pattern has a labyrinth structure.

In our numerical simulations, we apply a simple two-dimensional Ising-like model (see refs. 4, 5, and references therein). We consider a scalar spin field $\phi(\mathbf{r})$, where $\mathbf{r} = (x, y)$. Namely, the model is a continuous model. The positive and negative values of $\phi(\mathbf{r})$ correspond to the

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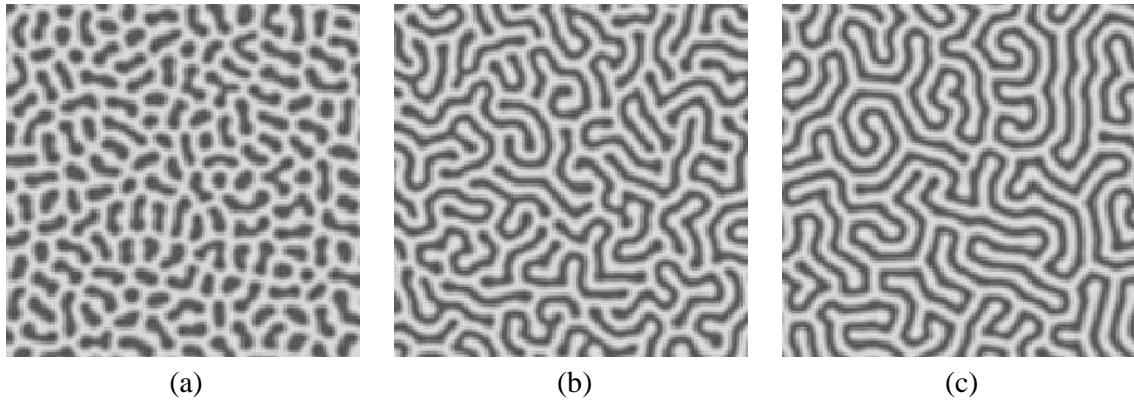


Fig. 2. Numerically simulated domain patterns under a zero field (1/4-parts of the whole system). The value of $\phi(\mathbf{r})$ is positive and negative in the white and black areas, respectively. The decreasing rate of the external field is (a) $v = 10^{-2}$, (b) $v = 10^{-3}$, and (c) $v = 10^{-4}$.

up and down spins, respectively. Our model Hamiltonian consists of four energy terms: uniaxial-anisotropy energy H_{ani} , local ferromagnetic interactions (exchange interactions) H_J , long range dipolar interactions H_{di} , and interactions with the external field H_{ex} . The anisotropy energy term implies that $\phi(\mathbf{r})$ tends to take either one of two different values [$\phi(\mathbf{r}) = \pm 1$]:

$$H_{\text{ani}} = \alpha \lambda(\mathbf{r}) \int d\mathbf{r} \left(-\frac{\phi(\mathbf{r})^2}{2} + \frac{\phi(\mathbf{r})^4}{4} \right), \quad (1)$$

where the factor $\lambda(\mathbf{r})$ expresses disorder, and

$$\lambda(\mathbf{r}) = 1 + \mu(\mathbf{r})/4. \quad (2)$$

Here, $\mu(\mathbf{r})$ is an uncorrelated random number with a Gaussian distribution whose average and variance are 0 and μ_0^2 , respectively. However, $\mu(\mathbf{r})$ should have a cutoff so that $\lambda(\mathbf{r})$ is always positive. We consider the disorder effect only in the anisotropy term, although some effects of disorder could also be applied to the other interaction terms. The ferromagnetic (exchange) and dipolar interactions are described by

$$H_J = \beta \int d\mathbf{r} \frac{|\nabla \phi(\mathbf{r})|^2}{2} \quad (3)$$

and

$$H_{\text{di}} = \gamma \int d\mathbf{r} d\mathbf{r}' \phi(\mathbf{r}) \phi(\mathbf{r}') G(\mathbf{r}, \mathbf{r}'), \quad (4)$$

respectively. Here, $G(\mathbf{r}, \mathbf{r}') \sim |\mathbf{r} - \mathbf{r}'|^{-3}$ at long distances. The integral in eq. (4) should have a lower cutoff since $G(\mathbf{r}, \mathbf{r}') \rightarrow \infty$ at short distances. The term from the interactions with the external field is given by

$$H_{\text{ex}} = -h(t) \int d\mathbf{r} \phi(\mathbf{r}), \quad (5)$$

where $h(t)$ is the time-dependent external field. From eqs. (1)–(5), the dynamical equation of our model is described by

$$\begin{aligned} \frac{\partial \phi(\mathbf{r})}{\partial t} &= -\frac{\delta(H_{\text{ani}} + H_J + H_{\text{di}} + H_{\text{ex}})}{\delta \phi(\mathbf{r})} \\ &= \alpha \lambda(\mathbf{r}) [\phi(\mathbf{r}) - \phi(\mathbf{r})^3] + \beta \nabla^2 \phi(\mathbf{r}) \\ &\quad - \gamma \int d\mathbf{r}' \phi(\mathbf{r}') G(\mathbf{r}, \mathbf{r}') + h(t). \end{aligned} \quad (6)$$

The coefficients of the anisotropy energy and the exchange and dipolar interactions — α , β , and γ respectively — are all positive.

In numerical simulations, it is useful to calculate the time evolution of eq. (6) in Fourier space. The equation is rewritten as

$$\frac{\partial \phi_{\mathbf{k}}}{\partial t} = \alpha [\lambda(\phi - \phi^3)]_{\mathbf{k}} - (\beta k^2 + \gamma G_{\mathbf{k}}) \phi_{\mathbf{k}} + h(t) \delta_{\mathbf{k}, 0}, \quad (7)$$

where $[\cdot]_{\mathbf{k}}$ denotes the convolution sum and $G_{\mathbf{k}}$ is the Fourier transform of $G(\mathbf{r}, 0)$. In the calculation below, if $G(\mathbf{r}, 0) \equiv 1/|\mathbf{r}|^3$, then

$$G_{\mathbf{k}} = a_0 - a_1 k, \quad (8)$$

where $k = |\mathbf{k}|$ and

$$a_0 = 2\pi \int_d^{\infty} r dr G(r), \quad a_1 = 2\pi. \quad (9)$$

Here, d is the cutoff length, which is approximately the thickness of the film in experiments. We set $d = \pi/2$, namely, $a_0 = 4$, in our simulations. In the simulations, 512×512 lattices with periodic boundary conditions are used, and we fix two parameter values: $\beta = 2.0$, $\gamma = 2\beta/a_1 = 2.0/\pi$. These parameters determine the characteristic length of domain patterns. According to the linear analysis of eq. (7), $\phi_{\mathbf{k}}$ of the modes with $k = 1.0$ grows the best at these fixed values. Namely, the characteristic wave number k_0 of the domain patterns is unity: $k_0 = 1.0$. Moreover, we fix the other two parameters as follows: $\alpha = 2.5$ and $\mu_0 = 0.3$. It would be interesting to investigate the dependence on α or μ_0 , which will be studied elsewhere. Only the changing rate v of $h(t)$ is the control parameter of our interest:

$$h(t) = h_{\text{ini}} - vt, \quad (10)$$

where h_{ini} is the saturating field; $h_{\text{ini}} \simeq 1.4$ in this case. The value h_{ini} was estimated from a simulation under an ascending field. In other words, the domains first spontaneously form a pattern under a zero field, and then an increasing field is applied. At a certain value of h , the domains with negative $\phi(\mathbf{r})$ disappear, and we define this value as h_{ini} .

Figure 2 shows the domain patterns under a zero field that are simulated by the abovementioned method. The white and black areas correspond to positive and negative $\phi(\mathbf{r})$, i.e. up and down spins, respectively. As the initial condition, $\phi(\mathbf{r})$ was given by random numbers in the interval $1.0 < \phi(\mathbf{r}) < 1.1$ at $t = 0$. The external field decreases from h_{ini} to 0, following eq. (10). Once the external field becomes zero at $t_0 = h_{\text{ini}}/v$, the field remains zero for $t \geq t_0$. The change in the domain patterns for $t \geq t_0$ is very little in this case. For some other sets of parameters (α and μ_0), however, the characteristics of the domain patterns change after the field vanishes.¹⁰

The simulation results have the same behavior as the experimental ones. Namely, a sea-island structure, which consists of many short domains, appears under a rapidly decreasing field, and a relatively small number of domains grow to form a labyrinth structure under a slowly decreasing field. In other words, both the experimental and numerical results suggest that the field sweeping rate affects the domain size (domain length).

The abovementioned results also mean that a domain pattern forms an unstable structure under a rapidly decreasing field. Since the ferromagnetic (exchange) interactions prefer small domain surfaces (domain walls), a structure composed of a few large domains is more stable than one composed of many small domains. In other words, a labyrinth structure is more stable than a sea-island structure from the viewpoint of energy. Before discussing why such an unstable structure appears under a rapidly decreasing field, we consider how domains grow to form a structure under a descending field.

How are the domain patterns in Fig. 2 created? The field dependence, which is now identified with the time dependence due to eq. (10), of the number and area of negative- ϕ (black) domains is shown in Fig. 3. Under a descending field, the total number of domains for each v rapidly increases at a certain field and has a plateau in the low-field region. In the plateau region, each domain grows; however, no new domain appears. The length of the plateau is longer for small v than for large v . For the rapidly decreasing field ($v = 10^{-2}$), the value of the field where the first domain appears is lower than that for the other slowly decreasing fields ($v = 10^{-3}, 10^{-4}$). From these results, we see that a small number of domains grow for a long time under a slowly decreasing field and vice versa.

Both the average and total domain areas show rapid growth at first and slow growth in the low-field region. This behavior is explained as follows. In the rapid growth region, where the domain area increases rapidly, domains grow freely. In the slow growth region, however, domains cannot grow freely since they tend to repel each other because of the dipolar interactions. Moreover, domains do not connect with each other for the parameter combination used in this calculation.¹¹ In the small- v case, since the number of domains is small, the change in the growth rates becomes clear. Here, we note that the difference between the total domain area and $512^2/2$ in Fig. 3(c) corresponds to the magnetization. We see that there is a small difference between the total domain area under a zero field and $512^2/2$, which corresponds to the remanent

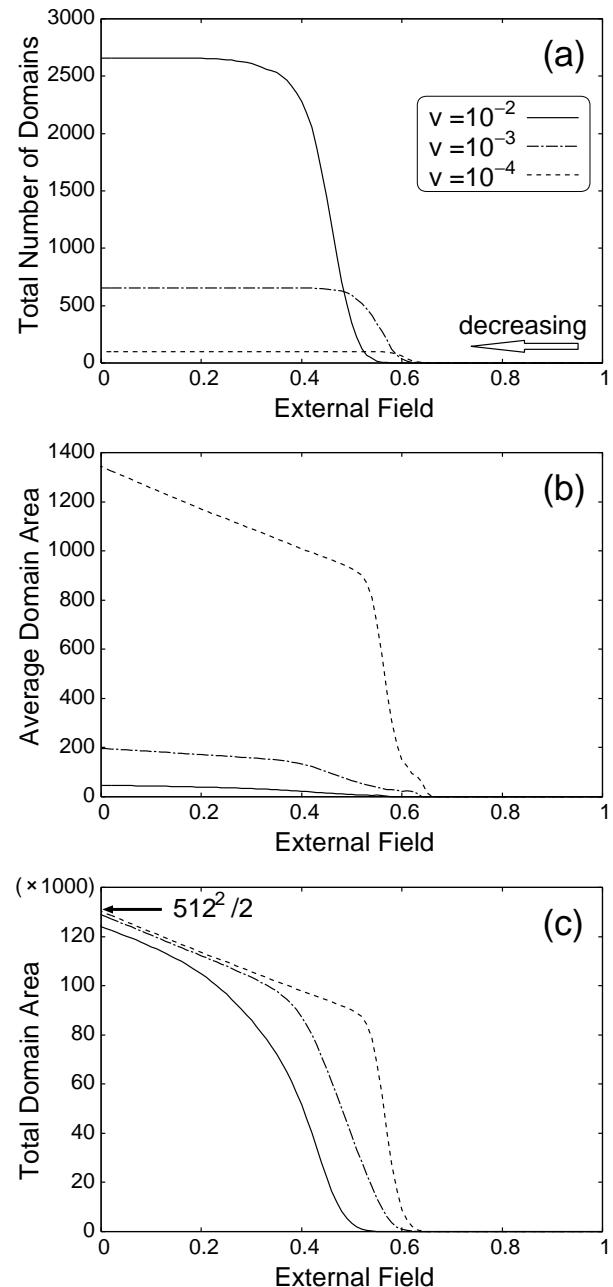


Fig. 3. The field dependence of (a) total number of domains, (b) average domain area, and (c) total domain area. The solid, chained, and dashed lines are for $v = 10^{-2}$, $v = 10^{-3}$, and $v = 10^{-4}$, respectively. The domain area is the number of grid points that have negative values of $\phi(\mathbf{r})$. The results under a zero field correspond to Fig. 2.

magnetization.

Now, let us discuss why the characteristics of domain patterns depend on the field sweeping rate, as shown in Figs. 2 and 3. We can explain the mechanism of this behavior by using the concept of crystallization: The driving force which causes nucleation corresponds to supersaturation or supercooling, and the nucleation energy ΔF is small for high supersaturation (supercooling).¹² In this system, the nucleation energy ΔF is the energy that is needed to create a negative- ϕ (black) domain in the "sea" of positive- ϕ spins (white space). However, there is

no defined quantity called supersaturation or supercooling. Let us consider a quantity that expresses the degree of non-equilibrium as “supersaturation.” The supersaturation should depend on the external field and the field sweeping rate. Since the changing rate of the valuable $\phi(\mathbf{r})$ is finite, the system cannot always keep up with the change in the field. It is natural to expect the following: the larger the value of v , the higher is the supersaturation. Therefore, when v is large, ΔF is small because of high supersaturation and nucleation easily occurs, which leads to the sudden appearance of many nuclei. On the contrary, when v is small, only a small number of nuclei appear because of large ΔF . Once the negative- ϕ domains appear, the domains grow but new domains hardly appear.

The supersaturation as well as the nucleation energy ΔF should be related to the hysteresis loop. In ref. 9, it is shown that the shape of the hysteresis loop changes depending on the sweeping rate of the field. The area of the hysteresis loop is equal to the energy dissipated per cycle. The area is large for a fast sweep and small for a slow sweep.⁹ Though the hysteresis and dynamic phase transition were mainly studied for different models in refs. 7–9, the methods used in the references will be helpful to evaluate ΔF . The quantization of ΔF will be discussed sometime in the future.

There is another question to be considered: how is it that a labyrinth structure is observed in some samples even when the external field is suddenly removed? This can be explained by the difference in the proper time scale of the individual samples in the absence of the applied field. The proper time scale is determined by the anisotropy, interactions between spins, and disorder in the samples. The difference in the time scale may be interpreted as the difference in “the sensitivity to the dynamical change in the field.” Using the concept of the *sensitivity* to the field change, we can consistently understand both our results and the above fact that a labyrinth structure often appears under a sudden decay of the field. Our results suggest that the *sensitivity*, in the same sample, should depend on the field sweeping rate v . On the other hand, if a sample has very high *sensitivity* to the field change, a labyrinth structure can be observed in the sample even when the external field is suddenly removed. In other words, the balance between the *sensitivity* and

the field sweeping rate is one of the factors that causes a variety of domain patterns.

In conclusion, we have confirmed that the characteristics of the magnetic domain patterns under a zero field depend on the decreasing rate of the external field in both the experiments and the numerical simulations. For a rapidly decreasing field, many short domains form a sea-island structure. On the contrary, a small number of long domains form a complex labyrinth structure for a slowly decreasing field. Those properties are explained qualitatively by using the idea of crystallization. In other words, in the case of large v , high “supersaturation” lowers the nucleation energy, which causes a sudden appearance of many nuclei, i.e., negative- ϕ domains, in a sea of positive- ϕ spins. However, a labyrinth structure appears even when the external field is suddenly removed in some experiments because the spins in actual garnet films usually vary fast enough to keep up with the change in the field. The results in this letter show that the characteristics of domain patterns are dependent on the field sweeping rate, or more generally, on the balance between the rate and a dynamical property such as “the sensitivity to the dynamical change in the field.”

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- 1) M. Seul and D. Andelman: *Science* **267** (1995) 476.
- 2) S. Miura, M. Mino and H. Yamazaki: *J. Phys. Soc. Jpn.* **70** (2001) 2821.
- 3) M. Mino, S. Miura, K. Dohi and H. Yamazaki: *J. Magn. Magn. Mater.* **226-230** (2001) 1530.
- 4) E. A. Jagla: *Phys. Rev. E* **70** (2004) 046204.
- 5) E. A. Jagla: *Phys. Rev. B* **72** (2005) 094406.
- 6) J. M. Deutsch and T. Mai: *Phys. Rev. E* **72** (2005) 016115.
- 7) C. N. Luse and A. Zangwill: *Phys. Rev. E* **50** (1994) 224.
- 8) S. W. Sides, P. A. Rikvold and M. A. Novotny: *Phys. Rev. E* **59** (1999) 2710.
- 9) H. Jang and M. J. Grimson: *Phys. Rev. E* **63** (2001) 066119.
- 10) We have observed such behavior in numerical simulations. The observations will appear in the future work.
- 11) For some other sets of parameters, the neighboring domains often connect with each other.
- 12) For example, I. Sunagawa: *Crystals* (Cambridge University Press, Cambridge, 2005), Chap. 3; See also other standard textbooks about crystal growth.